# Estimating the critical determinants of a class of three-dimensional star bodies

Werner Georg Nowak

**Abstract.** In the problem of (simultaneous) Diophantine approximation in  $\mathbb{R}^3$  (in the spirit of Hurwitz's theorem), lower bounds for the critical determinant of the special three-dimensional body

$$K_2: (y^2 + z^2)(x^2 + y^2 + z^2) \le 1$$

play an important role; see [1], [6]. This article deals with estimates from below for the critical determinant  $\Delta(K_c)$  of more general star bodies

 $K_c: (y^2 + z^2)^{c/2}(x^2 + y^2 + z^2) \le 1,$ 

where c is any positive constant. These are obtained by inscribing into  $K_c$  either a double cone, or an ellipsoid, or a double paraboloid, depending on the size of c.

## 1 Introduction

During the last couple of decades, not much research has been done in the subfield of the *Geometry of Numbers* (see, e.g., the monograph by Gruber & Lekkerkerker [4]) which is concerned with the evaluation, or at least estimation, of *critical determi*nants  $\Delta(K)$  of starbodies K in  $\mathbb{R}^s$ ,  $s \ge 2$ . These are defined as  $\Delta(K) = \inf |\det A|$ , where A ranges over all nonsingular real  $(s \times s)$ -matrices, such that the origin is the only point of the lattice  $A\mathbb{Z}^s$  in the interior of K.

It is the author's aim to rouse new interest in this classic topic *a fortiori* in view of its close connection to *simultaneous Diophantine approximation* in the spirit of Hurwitz's theorem: This is discussed at length in the author's survey article [9], as well as in the author's papers [6], [7], [8], [10].

<sup>2010</sup> MSC: 11J13, 11H16

 $K\!ey$  words: Geometry of numbers; critical determinant; simultaneous Diophantine approximation

DOI: 10.1515/cm-2017-0012

In brief, for each positive integer  $s \geq 2$ , and  $1 \leq \nu \leq \infty$ , define  $\theta_{s,\nu}$  as the supremum of all values C with the following property: For every  $\alpha \in \mathbb{R}^s \setminus \mathbb{Q}^s$ , there exist infinitely many  $(\mathbf{p}, q) \in \mathbb{Z}^s \times \mathbb{Z}_+$  with  $\operatorname{gcd}(\mathbf{p}, q) = 1$ , such that

$$\left\|\alpha - \frac{1}{q}\mathbf{p}\right\|_{\nu} < \frac{1}{q(Cq)^{1/s}}.$$
(1)

Then it is known due to a famous result of Davenport [3] that

$$\theta_{s,\nu} = \Delta(K^{(s,\nu)}), \qquad (2)$$

where

$$K^{(s,\nu)} = \left\{ (x_0, \dots, x_s) \in \mathbb{R}^{s+1} : |x_0| \, \| (x_1, \dots, x_s) \|_{\nu}^s \le 1 \right\} \,.$$

However, the exact determination of  $\theta_{s,\nu}$  has only been accomplished for s = 1 (Hurwitz's classic theorem:  $\theta_{1,\nu} = \sqrt{5}$ ) and for  $s = \nu = 2$ :  $\theta_{2,2} = \frac{1}{2}\sqrt{23}$  [2].

### 2 Objective of the present article

To fix notions, we concentrate on the most natural one of the unsolved cases concerning  $\theta_{s,\nu}$ , namely on our familiar three-dimensional space and the Euclidean norm. Armitage [1] proved that

$$\theta_{3,2} = \Delta(K^{(3,2)}) \ge (\Delta(K^*))^3 \Delta(K_2),$$

where

$$K^*: \quad x^2(x^2 + y^2)^3 \le 1$$

is a planar star body with  $\Delta(K^*) \ge 1.159$ , and

$$K_2:$$
  $(y^2 + z^2)(x^2 + y^2 + z^2) \le 1.$  (3)

Armitage proceeded to estimate  $\Delta(K_2)$  by inscribing an ellipsoid<sup>1</sup>

$$x^2 + 4y^2 + 4z^2 \le 2\sqrt{3} \,.$$

Thus he obtained

$$\theta_{3,2} = \Delta(K^{(3,2)}) \ge 1.774\dots$$

Around the turn of the millennium, the present author [6] replaced this ellipsoid by the double paraboloid

$$|x| \le (1 + \sqrt{2}) (1 - y^2 - z^2)$$
,

and evaluated the critical determinant of the latter. This gave the overall improvement

$$\theta_{3,2} = \Delta(K^{(3,2)}) \ge 1.879\dots$$

<sup>&</sup>lt;sup>1</sup>As it is common in the Geometry of Numbers, we will throughout use the terms *ellipsoid*, *paraboloid*, *cone*,  $\ldots$  for bodies, not for the boundary surfaces.

It is the aim of the present article to view the body  $K_2$  as a member of a more general family of star-bodies<sup>2</sup>

$$K_c: (y^2 + z^2)^{c/2} (x^2 + y^2 + z^2) \le 1,$$
 (4)

where c is an arbitrary fixed positive constant. Our objective is to deduce a lower bound for  $\Delta(K_c)$ , depending on c, for every c > 0.

We start with a brief survey of the bounds established, postponing a more detailed representation of the results to Table 2 at the end.

<i>c</i>	1	1.2	1.4	1.6	1.8	2
$\Delta(K_c) \ge$	0.9186	0.9612	1.0130	1.0780	1.1428	1.2071
С	2.2	2.4	2.6	2.8	3	
$\Delta(K_c) \ge$	1.2712	1.3358	1.4139	1.4917	1.5693	

Table 1: Lower bounds for  $\Delta(K_c)$  obtained, for a couple of values c.

#### 3 Strategy of proof and auxiliary results

There is no direct approach to estimate the critical determinant of a non-convex unbounded starbody like  $K_c$ . However, for convex (and **o**-symmetric) bodies in  $\mathbb{R}^3$ the situation is considerably better. For this case, Minkowski [5] has established a general theorem which tells us how in this case the critical lattices<sup>3</sup> necessarily look like; see also [4, p. 342, Theorem 3]. On the basis of this result, Minkowski was able to evaluate  $\Delta(\mathcal{O}) = \frac{19}{108}$  for the octahedron

$$\mathcal{O}: |x| + |y| + |z| \le 1.$$

Similarly, Ollerenshaw [11] showed that

$$\Delta(\mathcal{B}_3) = \frac{1}{\sqrt{2}} \tag{5}$$

for the origin-centered unit ball  $\mathcal{B}_3$  in  $\mathbb{R}^3$ . Furthermore, Whitworth [12] considered the double cone

$$\mathcal{C}: \quad |x| + \sqrt{y^2 + z^2} \le 1$$

and obtained

$$\Delta(\mathcal{C}) = \frac{\sqrt{6}}{8} \,. \tag{6}$$

Finally, the author [6] was able to show for the double paraboloid

 $\mathcal{P}: \quad |x| + y^2 + z^2 \le 1$ 

 $<sup>^{2}</sup>$ Obviously, no loss of generality is implied by the fact that only one of the exponents of the two brackets is assumed to vary.

<sup>&</sup>lt;sup>3</sup>A lattice  $A\mathbb{Z}^3$  is called *critical* for a body *B* if  $|\det A| = \Delta(B)$  and **o** is the only lattice point in the interior of *B*.

that

$$\Delta(\mathcal{P}) = \frac{1}{2}.\tag{7}$$

Our argument will be based on the idea to inscribe into  $K_c$  one of the three lastmentioned convex bodies, depending on the value of c, and to use the results (5)–(7). In fact, for a certain interval around c = 2, the choice of a paraboloid will turn out to be optimal, while for smaller values of c an ellipsoid will be the best choice, and for larger c the double cone will be most appropriate.

#### 4 The details of the analysis

**Lemma 1.** For fixed c, 0 < c < 4, let

$$\lambda_0 := \left(\frac{6}{4-c}\right)^{2/(2+c)} . \tag{8}$$

For any  $\lambda > 0$ , the ellipsoid

$$\mathcal{E}_{c}(\lambda): \quad \frac{x^{2}}{(1+\frac{1}{2}c)\lambda^{c/2}} + \frac{(y^{2}+z^{2})}{1+\frac{1}{2}c}\left(\frac{1}{2}c\lambda + \lambda^{-c/2}\right) \leq 1$$

is completely contained in  $K_c$  and has critical determinant

$$\Delta(\mathcal{E}_c(\lambda)) = \frac{(1 + \frac{1}{2}c)^{3/2}}{\sqrt{2}} \frac{\lambda^{c/4}}{\frac{1}{2}c\lambda + \lambda^{-c/2}}.$$
(9)

For any fixed c, 0 < c < 4, this expression attains its maximum for  $\lambda = \lambda_0$ , as given in (8). Hence  $\Delta(K_c) \ge \Delta(\mathcal{E}_c)$  with  $\mathcal{E}_c := \mathcal{E}_c(\lambda_0)$ .

*Proof.* Let  $r = \sqrt{y^2 + z^2}$  for short. Then, by the mean inequality with weights,

$$r^{c}(r^{2} + x^{2}) = (\lambda r^{2})^{c/2} \frac{r^{2} + x^{2}}{\lambda^{c/2}} \le \left(\frac{\frac{1}{2}c\lambda r^{2} + (r^{2} + x^{2})\lambda^{-c/2}}{1 + \frac{1}{2}c}\right)^{1 + c/2}$$

From this,  $K_c \supset \mathcal{E}_c(\lambda)$  is immediate. By (5), and an obvious linear substitution, (9) readily follows. Differentiating the right hand side of (9) with respect to  $\lambda$  and equating to zero, the choice  $\lambda = \lambda_0$ , as given in (8), turns out to be optimal.  $\Box$ 

**Lemma 2.** For any c > 0, define  $r_0(c)$  as the unique<sup>4</sup> solution in (0,1) of the equation

$$2r_0^{c+2} - (c+2)r_0 + c = 0.$$
<sup>(10)</sup>

Put further

$$x_0(c) := \frac{\sqrt{r_0(c)^{-c} - r_0(c)^2}}{1 - r_0(c)} \,. \tag{11}$$

152

<sup>&</sup>lt;sup>4</sup>Let L denote the left-hand side of (10), then  $\frac{dL}{dr_0} = 0$  iff  $r_0 = r_* := 2^{-1/(c+1)}$ . Hence L decreases on  $[0, r_*]$  from c to  $-r_*(c+1)$ , and increases on  $[r_*, 1]$  from  $-r_*(c+1)$  to 0. Hence the uniqueness of  $r_0(c)$ .

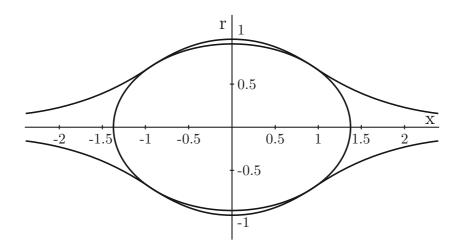


Figure 1: c = 1: The body  $K_1$  and an optimal inscribed ellipsoid, in front view

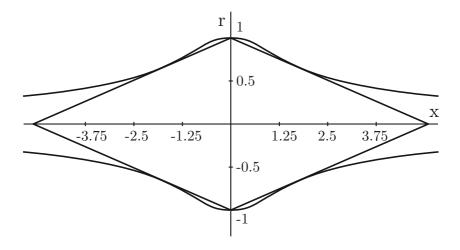


Figure 2: c = 3: The body  $K_3$  and an optimal inscribed double cone, in front view

Then the double cone

$$C_c: \quad \frac{|x|}{x_0(c)} + \sqrt{y^2 + z^2} \le 1$$

is completely contained in  $K_c$  and has critical determinant

$$\Delta(\mathcal{C}_c) = \frac{\sqrt{6}}{8} x_0(c) \,. \tag{12}$$

*Proof.* Since both  $K_c$  and  $C_c$  are bodies of rotation, with respect to the *x*-axis, it suffices to discuss the situation in front view - in a (x, r)-plane, say,  $r = \sqrt{y^2 + z^2}$ .

By symmetry, we may restrict the calculations to  $x \ge 0, r \ge 0$ . The curve  $k_c$  whose rotation generates  $\partial K_c$  is given by  $r^c(x^2 + r^2) = 1$ . Solving for x gives

$$x = \xi(r) := \sqrt{r^{-c} - r^2}$$
.

 $\partial \mathcal{C}_c$  is generated by the tangent

$$T: \quad x - \xi(r_0) = \xi'(r_0)(r - r_0) \tag{13}$$

which contains the point (x, r) = (0, 1). Inserting this into (13) and carrying out some bulky analysis, we arrive at (10). Since  $(x_0, 0)$  is the point of intersection of T with the x-axis, (11) follows by one more routine calculation. Finally, (6) readily implies (12).

By the way, the point of inflection of  $k_c$  is  $(\xi(r_W(c)), r_W(c))$  with

$$r_W(c) = \left(\frac{c}{2(c+1)}\right)^{1/(c+2)}$$

It is easily checked that throughout  $r_W(c) > r_0(c)$ .

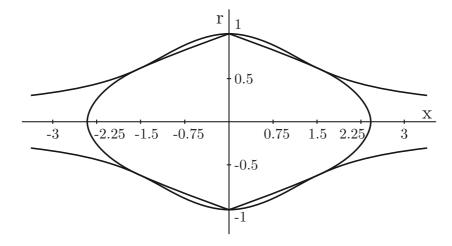


Figure 3: c = 2: The body  $K_2$  and an optimal inscribed double paraboloid, in front view

**Lemma 3.** For any c > 0, define  $r_1(c)$  as the unique<sup>5</sup> solution in (0,1) of the equation

$$2r_1^{c+2}(r_1^2+1) + c(1-r_1^2) - 4r_1^2 = 0.$$
(14)

Put further

$$\alpha(c) := \frac{\sqrt{r_1(c)^{-c} - r_1(c)^2}}{1 - r_1(c)^2} \,. \tag{15}$$

<sup>&</sup>lt;sup>5</sup>A similar argument applies as in footnote 4 in Lemma 2.

Then the double paraboloid

$$\mathcal{P}_c: \quad |x| \le \alpha(c)(1 - y^2 - z^2)$$

is completely contained in  $K_c$  and has critical determinant

$$\Delta(\mathcal{P}_c) = \frac{\alpha(c)}{2} \,. \tag{16}$$

Proof. Again we consider the situation in front view, in (x, r)-variables,  $x, r \ge 0$ . The aim is to choose  $\alpha = \alpha(c)$  so that the parabola  $p_c$ :  $x = \alpha(1 - r^2)$  and the curve  $k_c$  have one point  $(x_1, r_1)$  in common (in the first quadrant), where also the derivative  $x'_1 = \frac{dx}{dr}\Big|_{r=r_1}$  has the same value. In this way we get:

$$r_1^c(r_1^2 + x_1^2) = 1, (17)$$

$$(c+2)r_1^{c+1} + cr_1^{c-1}x_1^2 + 2r_1^c x_1 x_1' = 0, \qquad (18)$$

$$x_1 = \alpha (1 - r_1^2), \tag{19}$$

$$r_1' = -2\alpha r_1 \,. \tag{20}$$

Dividing (20) by (19), we conclude that

$$x_1' = -\frac{2r_1}{1 - r_1^2} \, x_1 \, .$$

Using this in (18), we get

$$(c+2)r_1^{c+1} + x_1^2 \left( cr_1^{c-1} - \frac{4r_1}{1-r_1^2} \right) = 0.$$
(21)

Solving (17) for  $x_1^2$  and using this in (21), we obtain an equation in the single unknown  $r_1$  which, after simplifying, is just (14). Further, (17) and (19) readily imply (15). Finally, (16) is immediate from (7).

Again, it is easily checked numerically that throughout  $r_W(c) > r_1(c)$ .

We are now in a position to summarize the results obtained.

**Theorem 1.** For 0 < c < 4, the critical determinant of the starbody

$$K_c: \quad (y^2 + z^2)^{c/2}(x^2 + y^2 + z^2) \le 1$$

can be estimated from below by

 $\Delta(K_c) \geq \max(\Delta(\mathcal{E}_c), \Delta(\mathcal{C}_c), \Delta(\mathcal{P}_c)).$ 

Here,  $\Delta(\mathcal{E}_c), \Delta(\mathcal{C}_c), \Delta(\mathcal{P}_c)$  are given in Lemmas 1–3. Further, for  $c \geq 4$ ,

$$\Delta(K_c) \ge \Delta(\mathcal{C}_c) \,.$$

**Remark 1.** As can be seen from the table below, for  $c \in \{1, 1.2\}$ , the sharpest lower bound for  $\Delta(K_c)$  can be obtained by inscribing an ellipsoid. For  $c \in \{1.4, 1.6, 1.8, 2, 2.2\}$ , inscribing a double paraboloid yields the best result, while for  $c \in \{2.4, 2.6, 2.8, 3\}$ , an inscribed double cone is the best choice.

c	$\Delta(\mathcal{E}_c)$	$\Delta(\mathcal{P}_c)$	$\Delta(\mathcal{C}_c)$
1	0.9186	0.8810	0.7785
1.2	0.9612	0.9473	0.8601
1.4	1.0045	1.0130	0.9408
1.6	1.0485	1.0780	1.0207
1.8	1.0935	1.1428	1.1000
2	1.1398	1.2071	1.1790
2.2	1.1875	1.2712	1.2576
2.4	1.2371	1.3351	1.3358
2.6	1.2890	1.3988	1.4139
2.8	1.3437	1.4623	1.4917
3	1.4019	1.5257	1.5693

Table 2: The critical determinants of  $\mathcal{E}_c, \mathcal{P}_c, \mathcal{C}_c$ , for  $1 \leq c \leq 3$ , in step lengths of 0.2.

#### References

- J.V. Armitage: <u>On a method of Mordell in the geometry of numbers</u>. Mathematika 2 (2) (1955) 132–140.
- [2] H. Davenport, K. Mahler: <u>Simultaneous Diophantine approximation</u>. Duke Math. J. 13 (1946) 105–111.
- [3] H. Davenport: On a theorem of Furtwängler. J. London Math.Soc. 30 (1955) 185–195.
- [4] P.M. Gruber, C.G. Lekkerkerker: Geometry of numbers. North Holland, Amsterdam (1987).
- [5] H. Minkowski: Dichteste gitterförmige Lagerung kongruenter Körper. Nachr. Kön. Ges. Wiss. Göttingen (1904) 311–355.
- [6] W.G. Nowak: The critical determinant of the double paraboloid and Diophantine approximation in R<sup>3</sup> and R<sup>4</sup>. Math. Pannonica 10 (1999) 111–122.
- [7] W.G. Nowak: Diophantine approximation in ℝ<sup>s</sup>: On a method of Mordell and Armitage. In: Algebraic number theory and Diophantine analysis. Proceedings of the conference held in Graz, Austria, August 30 to September 5, 1998, W. de Gruyter, Berlin. (2000) 339–349.
- [8] W.G. Nowak: Lower bounds for simultaneous Diophantine approximation constants. Comm. Math. 22 (1) (2014) 71–76.
- [9] W.G. Nowak: Simultaneous Diophantine approximation: Searching for analogues of Hurwitz's theorem. In: T.M. Rassias and P.M. Pardalos (eds.), Essays in mathematics and its applications. Springer, Switzerland (2016) 181–197.
- [10] W.G. Nowak: On the critical determinants of certain star bodies. Comm. Math. 25 (1) (2017) 5–11.
- [11] K. Ollerenshaw: The critical lattices of a sphere. J. London Math. Soc. 23 (1949) 297–299.
- [12] J.V. Whitworth: The critical lattices of the double cone. Proc. London Math. Soc. 2 (1) (1951) 422–443.

Author's address: INSTITUTE OF MATHEMATICS, DEPARTMENT OF INTEGRATIVE BIOLOGY, BOKU WIEN, 1180 VIENNA, AUSTRIA

 $E\text{-}mail: \verb"nowak@boku.ac.at"$ 

Received: 21 February, 2017 Accepted for publication: 6 July, 2017 Communicated by: Karl Dilcher